Particle ID for top, Higgs and SUSY

oTau ID (end of it) oJets and Missing E_{T} , neutrinos oTrigger strategy and rates oW and Zs oB tagging oTop physics oHiggs search oSUSY oConclusions

Offline Tau identification in Atlas



At 50 GeV (green) QCD jet rejection >1000 for τ -jet acceptance=50% 2

Jets

- · Jets are comparatively easier to trigger on and reconstruct.
- Cross-section decreases very fast with E_T accurate E_T measurement at trigger level is important
 - \rightarrow large cluster size like 0.8 x 0.8 or more
 - \rightarrow correct weighting of EM and HCAL energies (ATLAS and CMS

calorimeters are non-compensating ...)

• Ability to separate nearby jets \rightarrow smaller cluster size preferred

ATLAS works with a 4x4 window of 0.2x0.2trigger cells
A LVL1 internal logic eliminates dble counting and finds core of triggering jet, which defines RoI for HLTall that every 25ns for the whole solid angle...





Figure 4-31 Jet trigger efficiency curves for 100 GeV $E_{\rm T}$ jets, for different cluster sizes, at luminosity 10^{33} cm⁻²s⁻¹.

Figure 4-32 Trigger rate vs. efficiency for 100 GeV E_T jets, for different cluster sizes, at luminosity 10³³ cm⁻²s⁻¹.

Missing E_T at the trigger/reconstruction level

•From the position and energy of each of the calorimeter cells, are calculated, summing on EM and HCAL sections.

- ΣE_x and ΣE_y a 2-vector in the transverse plane whose modulus is E_T miss - ΣE_T in the transverse plane, also called "total E_T "

•If there are no missing particles $\Sigma E_x = 0$ and $\Sigma E_y = 0$, ie MET= 0 •Accuracy limited by :

-fluctuations of sampled energies, and noise (option=threshold) -uncovered solid angle (η >5),(high E, but *sin(θ) \rightarrow 0=OK)

-cracks,...

 Conversely E_T miss >Eth signs (a) missing particle(s): a neutrino(s) or something more exotic....

LVL1: use calo trigger cells HLT and offline: use calo cells, apply weighting and zero suppression for better MET resolution



From missing E_T to missing particle(s)

Need hypotheses....to be confirmed by event analysis: •<u>Single particle missing</u> (v,neutralino,..) MET = its transverse momentum

•<u>Two particles missing</u> = ambiguous in the transverse plane. can be solved (transverse and longitudinal) if missing particles are decay products of two "massless" parents, like taus, of which other decay particules are identified (as a narrow jet).



Missing E_T in the trigger....

LVL1 MET trigger

 \rightarrow too high rate in stand alone to catch for example $W \rightarrow \tau v$ \rightarrow use it combined with other signatures: $-E_{T}$ miss +taus $-E_{T}$ miss +jets (SUSY),....





Figure 15-47 Event rates as function of E_T^{miss} when requiring a jet above various thresholds. Left: low luminosity; right: high luminosity.

Expected LVL1 rates at "low" L

Trigger CMS 10 ³³	Threshold (GeV or GeV/c)	Rate (kHz)	Cumulative (kHz)	RateATLAS
Inclusive isolated electron/photon	29	3.3	3.3	20GeV/11 kHz
Di-electrons/di-photons	17	1.3	4.3	15GeV/2 kHz
Inclusive isolated muon	14	2.7	7.0	6GeV/23 kHz
Di-muons	3	0.9	7.9	
Single tau-jet trigger	86	2.2	10.1	
Two tau-jets	59	1.0	10.9	20-30/2kHz
1-jet, 3-jets, 4-jets	177, 86, 70	3.0	12.5	180-75-55/0.6
$Jet * E_T^{miss}$	88 * 46	2.3	14.3	50*50/0.4 kHz
Electron * Jet	21 * 45	0.8	15.1	
Minimum-bias (calibration)		0.9	16.0	
TOTAL			16.0	40 kHz

Table 15-1 Level-1 Trigger table at low luminosity. Thresholds correspond to values with 95% efficiency.

HLT reduces to <~200 Hz the rate to "permanent storage", keeping the thresholds energies at or very close to the LVL1

A possible strategy

Selection signature	Examples of physics coverage	Examples of physics coverage			
e25i	$W \rightarrow ev, Z \rightarrow ee, top production, H \rightarrow WW^{(*)}/ZZ^{(*)}, W',Z'$				
2e15i	$Z \rightarrow ee, H \rightarrow WW^{(*)}/ZZ^{(*)}$				
μ20i	$W \to \mu \nu, Z \to \mu \mu,$ top production, $H \to WW^{(*)}/ZZ^{(*)}, W^{*}, Z^{*}$				
2μ10	$Z \rightarrow \mu\mu$, $H \rightarrow WW^{(*)}/ZZ^{(*)}$				
γ60i	direct photon production, $H\to\gamma\gamma$				
2γ201	$H \rightarrow \gamma \gamma$				
j400	QCD, SUSY, new resonances				
2j350	QCD, SUSY, new resonances	QCD, SUSY, new resonances			
3j165	QCD, SUSY	QCD, SUSY			
4j110	QCD, SUSY	ATLAS 2 x 10 ³³			
τ60i	charged Higgs	tinal selection			
µ10 + e15i	$H \rightarrow WW^{(*)}/ZZ^{(*)}$, SUSY	$H \rightarrow WW^{(*)}/ZZ^{(*)}$, SUSY			
τ35i + xE45	$qqH(\tau\tau),W\to\tau\nu,Z\to\tau\tau,SUS$	${\rm qq}H(\tau\tau),W\to\tau\nu,Z\to\tau\tau,SUSY$ at large tan β			
j70 + xE70	SUSY	SUSY			
xE200	new phenomena	new phenomena			
E1000	new phenomena	new phenomena			
jE1000	new phenomena				
$2\mu 6 + \mu^+\mu^- + mass cuts$	rare b-hadron decays (B $ ightarrow \mu\mu$)	() and $B \rightarrow J/\psi (\psi')X$			

W and Zs to calibrate the detector and make important SM measurements

From cross-section, acceptance (η < 2.5 and trigger) & luminosity \Rightarrow event rate Assuming 100 days at 2 10³³ gives:

-5 10° Z \rightarrow ee and 5 10° Z \rightarrow µµ to mass storage (0.5 Hz each)

-5 10⁷ W \rightarrow ev and 5 10⁷ W \rightarrow µv ""



Triple Gauge bosons couplings and probe SM.

Vertex detectors and B-tagging



Atlas pixels



Pixels electronics

FE electronics and readout architecture

-80 million channels to be "looked at" every 25 ns !!

-fastpreamp, good S(15ke)/N (200e) for mips \rightarrow digital or few bits r/o

-all FE chips proceed in parallel, controlled by local MicroControler

-occupancy small (<<10⁻²) : logic by columns

Schematically:

-each FE chip covers 24 columns of 160 cells [400 \times 50 μ m]=~1cm²

-at the bump bond pad is connected preamp + discriminator

-the End of Column logic drains data from columns=

address of hit pixel(s) in the column +bunch counter

-upon LVL1 signal(<100kHz) the MC scans over its FEchips(~12) gather the hit pixel adresses +bc ,keeps those with proper bc,and clears the buffers.





B-tagging expected performance



- •Start from a calorimeter jet (incl tau-jet..)
- •Reconstruct tracks, and select those with p_T >cut (~1 GeV/c) and ΔR <0.4
- •Measure d (with a sign), calculate s=d/ σ
- •Calculate jet weight as $\Sigma \log(\text{signif as b/signif as u})$ from all its tracks
- Adjust cut position
- •Positive tails in u-jets ? secondary interactions, Vo
- •Negative tails in b-jets? Jet direction \neq B direction

u-iet

b-jet

20

30

Jet weight

40

B-tagging expected performance



efficiency each (BR=10% each)

 $\varepsilon_{\rm h}$

B-tagging in HLT?

•CDF is using it for selection of "unbiased" B decays

Starting from LVL1 with 2 tracks >2 GeV/c pT
Done in a "hardware oriented" way with a processing time of 25 µs/event

- 12 independent sectors
- 4 layers(out of 5) of micro-strips grouped to 250 μm pitch
- 1 point/tracker (φ, p_T) + hit pattern compared to 32000 masks



•Interesting approach in LHCb/Velo

In ATLAS-CMS events to be processed, at LVL2 have in general several jets Some trial at LVL2-Atlas, not really convincing
Need "real tracking" to decide if some of the jets are B-jets or not. Efficient way to do it is at the Event Filter level, with ~full offline performance

B-tagging in HLT:CMS example



Figure 15-68 Efficiency for the *b*-tag versus mistagging rate for jet with E_T=100 GeV in the low (left) and high (right) luminosity scenarios.

- •Reconstruct Tracks in a cone defined by Calo-jet
- •Ask for at least 2/3 pixel hits.
- •Labelled as b-jet if ≥ 2 tracks above thresh significance 17

Particle ID for M(top) measurements

M(W), M(top) and M(Higgs) are linked in the standard model:

 $M(w) = K(\alpha, sin(\Theta_W), G_F) / sqrt(1 - \Delta r) + F[M(top), log(M(H)]]$

 $G_{\rm F}$, $a_{\rm EM}$, $\sin\theta_{\rm W}$ known with high precision, precise measurements of M(top) and M(W) constrain M(H) (weakly because of log term) 10 MeV M(W) $\rightarrow \sim 1$ GeV M(top) To day (CDF+D0) M(top) = 172.5 \pm 2.3 GeV Best strategy to measure it at LHC? [event statistics is not the main problem]



W

Particle ID for M(top) measurements



Inclusive tt cross-section at LHC~0.7 nb (100 x Tevatron)

t→bW 100%

 $W \rightarrow lv$ 11% each , 67% hadronic; no Bs

 $W1 \rightarrow had W2 \rightarrow had$ $W1 \rightarrow e/\mu v$, $W2 \rightarrow had$ (+ 2 \rightarrow 1) : 30% 1 mass fully reconstructed W1 \rightarrow e/µ,v ,W2 \rightarrow e/µv Others: one $W \rightarrow \tau$

: 44% 2 masses fully rec/ huge QCD BG : 5% rate; mass not fully reconstructed : not appropriate for precise M meast



Particle ID for M(top) measurement

Selection efficiency. = 5%Selection \Rightarrow 126k events, with S/B ~65 - 1 iso lepton, pT > 20 GeV, $|\eta| < 2.5$ - pTmiss > 20 GeV Events/4 GeV - \ge 4 jets with pT > 40 GeV, $|\eta| < 2.5$ $- \ge 2$ jets with b-tag 10000 p-1 > 20GeV As before, As before, Events per 10 fb⁻¹ E_miss > 20GeV plus N_{iet}≥4 Process plus N_{b-iet}≥2 5.0 126 000 tt signal 64.7 21.2 0 100 p_{T}^{200} (lepton) (GeV) q_{T}^{300} 0.002 1658 W+jets 47.9 0.1 15.0 0.05 0.002 Z+jets 232 Events 0000015 10000 0.5 0.006 53.6 10 ww 0.02 53.8 0.5 8 WZ 7500 2.8 0.04 0.008 14 ZZ 5000 1922 Total background Efficiency of cuts S/B 65 2500

¹⁰ Njets

0

W-ID from "top" sample

- •Semi-leptonic ttbar events contain W→jet-jet evts with good S/N ⇒identified hadronic W decays
- •This sample of jets is used to adjust the (non-b) jet -scale, starting from jets normalized using photon-jet evts (p_T balance)
- B-jets are normalized from photon(or Z)+ B-id jets (only)
- Remaining uncertainties on jet scales dominate the top mass systematic uncertainty

•Hadronic Ws in jj used in other places (H \rightarrow WW,..)



SM-Higgs search global view





The ZZ "gold plated" mode





HO

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Higgs mass (GeV)	200	240	280	320	360	400	500
$BR(H \rightarrow ZZ)$	0.26	0.29	0.30	0.31	0.30	0.28	0.27
$\sigma imes BR$ (fb)	12.4	11.2	9.6	8.9	8.7	6.8	3.2
Signal (no p _T cut)	134	127	110	105	105	86	44
Background (no p_{T} cut)	74	57	43	33	29	29	17
S / \sqrt{B} (no p_{T} cut) for 30 fb ⁻¹	15.6	16.8	16.8	18.2	19.3	15.9	10.7

-what about one Z in jet-jet? BG "Z +2 jets" is large....some attempt with one Z in two b jets ("pedagogical" reaction of the beginning) In fact, does not bring much....

ZZ* 130<M(H)<180 GeV

- still 4 leptons ee or $\mu\mu$

- only one M(II)=M(Z) constraint remaining
- dip in the BR at 2M(W)
- -Zbb main reducible background can be reduced applying a veto on displaced vertices -important issue : <u>acceptance of the lower PT</u> <u>lepton</u> .Analyses require 2 lepons pT>20 GeV (trigger) and 2 leptons with pT>7 GeV (offline)



fb⁻¹

Table 19-14 Signal and background rates after all cuts and signal significances as a function of m_{H^*} for $H \rightarrow ZZ^* \rightarrow 4I$ events and for an integrated luminosity of 30 fb⁻¹ (low luminosity performance).

Higgs mass (GeV)	120	130	150	170	180	
Signal	4.1	11.4	26.8	7.6	19.7	10
ŧī	0.01	0.02	0.03	0.02	0.02	
Zbb	0.08	0.12	0.19	0.17	D.19	
ZZ*	1.23	2.27	2.51	2.83	2.87	30
$ZZ \rightarrow \tau \tau II$	0.13	0.20	0.25	0.08	0.02	
Significance ($S \land \sqrt{B}$)	3.4	7.0	15.5	4.3	11.2	
Significance (Poisson)	2.4	4.8	15.5	3.2	11.2	

In 50% of cases the softer lepton has Pt < 10 GeV/c (not a trigger problem)



High luminosity

Higgs in 2 photons

For M(H)<150 GeV, most promising channels are: <u> $H \rightarrow \gamma\gamma$ (inclusive)</u> and <u> $H \rightarrow \tau\tau$ in association with 2 forward jets (VBF)</u>

Distinctive features of the $H \rightarrow \gamma \gamma$ mode are:

-limited $\sigma \ x \ BR$

- -large irreducible background
- -potentially large instrumental background (jet-jet and γ -jet)
- -clean signature -high invariant mass resolution

defining criteria for EM calorimeters

Ultimate photon energy resolution , when all calibration/normalization problems are solved: CMS : $3\%/\sqrt{E} \oplus (200 \text{MeV} \oplus \text{pile-up})/E \oplus 0.55\%$ (??) ATLAS : $10\%/\sqrt{E} \oplus (200 \text{MeV} \oplus \text{pile-up} / E \oplus 0.70\%$ (?) Remember converted photons are somewhat worse $\rightarrow \Delta M(H)/M(H) \sim 1\%$ is the (difficult) target

Jet-jet and γ -jet rejection

Overall jet rejection obtained in ATLAS MC full simulation, confirmed by test beam (using fine strips):

-1050 for quark jets

-6000 for gluon jets \rightarrow Ultimate performance process dependant! (probability of a high x isolated π^0 is much higher in (MC) quark jet



 10^{-3}

10⁻⁴L

Kniehl et al

0.1 0.2 0.3 0.4 0.5 0.6 0.7

0.8 0.9 1 momentum fraction

Channel very much dependant on detector Ultimate performance, resolution and photon/ jet identification.

Signal and background in CMS



Higgs $\rightarrow \tau\tau$ by VBF

The ττ mode has a larger BR than γγ but a weaker signature and worse resolution
 Asking 2 forward jets (and no other jet) selects weak process by WW fusion, which has less background





 $m_{ev}\left(GeV\right)$

VBF, jet tagging, jet veto



Forward region is a difficult one:

- •transition at η =3 between geometries and/or technologies
- •Jet tagging sensitive to pile-up noise
- ·Jet veto problematic at high luminosity

SUSY (just a glance..)

•In the MSSM each fermion $f_{L,R}$ has a scalar partner $sf_{L,R}$ and to each Gauge boson is associated a massless spin $\frac{1}{2}$ gaugino.

•Among other virtues, these states coming with an opposite sign to normal particles in loop corrections, cancel the quadratic divergence of the Higgs mass, which would otherwise run to "infinity".

•None of these particles has been so far observed \rightarrow heavy ! •sq and sg couple to QCD as usual particles, they should be copiously produced at LHC <u>as soon as threshold is passed.</u>

 In the MSSM R-parity = (-1)**3(B-L)+2S is conserved: s-particles are produced in pair the lightest one is stable= neutralino χ⁰

•Exact signatures depend on the details of the model

<u>A common feature is large ΣE_T and missing energy (neutralinos)</u>

• The reach of experiments can be evaluated and compared using as parameter m_0 ($m_{1/2}$) the common mass of all scalars (spin $\frac{1}{2}$) at GUT scale



SUSY parameters from decay chain



A,H decays in $\tau\tau$

•In the MSSM there are 5 Higgs bosons : h, H, A, H⁺⁻ The 5 masses depend on 2 parameters: M(A) and $tg(\beta)$

•For large M(A) and tg(β) the H and A couplings to bb and τ τ are enhanced bb—bbA

•LVL1 Triggers = (one tau-jet or lepton)⊕ETmiss



Conclusion-Summary

- Each of the 4 LHC experiment uses technologies at the limit, for ambitious physics program
- Particle identification plays a major role
- Detectors have been carefully designed and are being built, installed and tested with a lot of care.
- Tuning the triggers at start-up will be a crucial step
- A major issue-not discussed here-will be the capacity of the experiments to "digest" the enormous amount of data they will produce, typically 1 Mbyte*200Hz to mass storage
- More in 1.5 years, when LHC gets started

Back-up transparencies



MSSM	t	b/ au	W/Z
/SM			
h	cos A/	-sinα/	$sin(\alpha - \beta)$
	$sin\beta$	$cos\beta$	
H	sin A/	Cos X/	$cos(\alpha - \beta)$
	sineta	$cos\beta$	
A	$cot\beta$	taneta	

Light Higgs boson h: 300 fb⁻¹



An example of exotica :Black hole production..

•Another way of stabilizing the Higgs mass is to assume that there are extra space-dimensions, in which case the Planck mass could be as low as a few TeV.

 •When √S reaches Mp Black Holes are produced, which decay "democratically" to quarks, leptons, Gauge bosons.



A typical event (?) could have: hadrons/leptons/γ,W,Z/Higgs ~ 75%/20%/3%/2%

....a dream for particle identification at the LHC